# Loop Antenna Equations

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*Abstract*—A brief discussion of a loop antenna in empty space that also serves to catalog some important antenna relations. Revised: June 6, 2007

## I. LOOP RADIATION

An excellent discussion of loop antenna physics can be found in McDonald's webnotes [1]. Therein the general expressions for the radiation properties of any parallelogram loop are considered. Assuming a square loop a meters on a side we can calculate for the radiation intensity (in Watts/steradian) [2]

$$U(\theta,\phi) = \frac{2I_0^2}{\pi^2} \sqrt{\frac{\mu_0}{\epsilon_0}} \frac{\sin^2(\frac{\pi a}{\lambda}\sin\theta\cos\phi)\sin^2(\frac{\pi a}{\lambda}\sin\theta\sin\phi)}{\sin^2\theta\sin^22\phi}$$
(1)

where  $I_0$  is the amplitude of the sinusoidal current flowing in the loop,  $\lambda$  is the wavelength of the radiation signal,  $\theta$  is the elevation angle (with respect to the plane of the loop),  $\phi$  is the azimuth angle (in the plane of the loop),  $\mu_0$  is the permeability of vacuum, and  $\epsilon_0$  is the permittivity of vacuum.

The total amount of power absorbed/emitted by this antenna is

$$P_{rad} = \frac{I_0^2 R_r}{2} \tag{2}$$

where  $R_r$  is the radiation resistance. Integrating the reference power density over the spherical solid angle results in the total radiation power, that is

$$P_{rad} = \int_{\Omega} U(\theta, \phi) \mathrm{d}\Omega \tag{3}$$

so that

$$R_r = \frac{4}{\pi^2} \sqrt{\frac{\mu_0}{\epsilon_0}} \int_{\Omega} \frac{\sin^2(\frac{\pi a}{\lambda} \sin \theta \cos \phi) \sin^2(\frac{\pi a}{\lambda} \sin \theta \sin \phi)}{\sin^2 \theta \sin^2 2\phi} d\Omega.$$
(4)

For comparison the radiation resistance of a "small" loop antenna (i.e. dimensions  $\ll \lambda$ ) without ground plane is [3]

$$R_r = 3.12 \times 10^4 \left(\frac{NS}{\lambda^2}\right)^2 = 3.12 \times 10^4 \left(\frac{N(l_1 \cdot l_2)}{\lambda^2}\right)^2 \qquad [\Omega]$$
(5)

where N is the number of loops, S the loop area, and  $l_1$  and  $l_2$  the antenna dimensions in the case of a square loop. Ideally the antenna dimensions will conform to those of the integrated technology around which the loop is wrapped [4]. This means that we can expect  $l_1$  and  $l_2$  to range from 1 mm or less up to 15 mm — probably a rather aggressive value give that microprocessors typically range about 10 mm on a side. If we consider possible frequencies of operation up to say 60 GHz, we must resort to a more involved account of the antenna properties, as presented above, since the wavelength at these

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frequencies (5 mm at 60 GHz) is on the same order as the antenna dimensions.

If we know enough about the material details of the loop antenna we can calculate the loss resistance,  $R_l$ , and hence the antenna efficiency,

$$\eta = \frac{R_r}{R_r + R_l}.$$
(6)

This factor has a place in the consideration of the antennas effective aperture area (the aperture),  $A_e$ . We can think of  $A_e$  as the area through which the radiative flux power appearing at the antenna's terminals had to pass (to reach the antenna's terminals). From [2] we recall that

$$A_e(\theta,\phi) = \lambda^2 \frac{U(\theta,\phi)}{P_{inj}} = \eta \lambda^2 \frac{U(\theta,\phi)}{P_{rad}}$$
(7)

where  $P_{inj}$  is the power appearing at the antenna's terminals (i.e. either coming from the transmitter circuit in transmit mode or available to the receiver in receive mode). Clearly the lower the carrier frequency the greater the antenna's aperture.

Substituting Eq. (6) into Eq. (7) we get

$$A_e(\theta,\phi) = \frac{2\lambda^2}{R_r + R_l} \frac{U(\theta,\phi)}{I_0^2}.$$
(8)

Two other important antenna measures are the gain  $G(\theta, \phi)$ and directivity  $D(\theta, \phi)$ . The key relationships for these variables are

$$G(\theta,\phi) = \frac{4\pi}{\lambda^2} A_e(\theta,\phi) = \frac{4\pi U(\theta,\phi)}{P_{inj}}.$$
(9)

As noted in [2] the when referring to the gain without any arguments, that is, as just G, we are implicitly referring to the maximum antenna gain. The directivity is simply the gain corrected for the antenna efficiency

$$D(\theta, \phi) = \frac{G(\theta, \phi)}{\eta} = \frac{4\pi U(\theta, \phi)}{P_{rad}}.$$
 (10)

### **II. LOOP INDUCTANCE**

The inductance of a rectangular loop of wire of radius rand dimensions  $l_1 \times l_2$  is given by Stutzman [3] as

$$L = \frac{\mu_0}{\pi} \left( l_2 \cdot \operatorname{acosh} \frac{l_1}{2r} + l_1 \cdot \operatorname{acosh} \frac{l_2}{2r} \right)$$
(11)

Since we will be dealing with planar coil designs, a suitable approximation to r for a planar coil of line width, w, is

$$r \approx w/2\pi.$$
 (12)

The assumption is that we can treat the width of the planar coil as the circumference of circular coil.

#### **III. LOOP LOSSES**

In calculating the loop's resistive losses we start by noting that the skin depth of a metal is given by

$$\delta = \sqrt{\frac{2}{\omega\mu_0\sigma}} \tag{13}$$

where  $\omega_0$  is the (angular) frequency of the signal through the metal with conductivity  $\sigma$  [ $\Im/m$ ]. With this we can approximate the series loss of the coil with

$$R_l \approx \frac{2(l_1 + l_2)}{w \cdot \delta \cdot \sigma} \tag{14}$$

where  $1/\sigma$  is multiplied by the total length of the loop (numerator) and divided by the effective cross-sectional area of the loop. Lee [5] (pg. 141) refines this approximation further to

$$R_l \approx \frac{2(l_1 + l_2)}{w \cdot \delta(1 - e^{-t/\delta}) \cdot \sigma}.$$
(15)

As noted by Lee, in actual integrated scenarios, the substrate losses and proximity effects would have to be accounted for in an accurate calculation of  $R_l$ .

## **IV. INSIGHTS**

From Eq. (5) we can see that the radiation resistance of a small loop antenna is proportional to  $f^4$  while  $R_l$  (Eq. (15)) is proportional to  $\sqrt{f}$ , hence our expressions for  $R_r$  and  $R_l$  combine to predict a rather poor antenna efficiency,  $\eta$ , at low carrier frequencies. Example calculations are given in [4].

#### REFERENCES

- [1] K. T. McDonald, "A parallelogram loop antenna," May 28 2003, URL: http://physics.princeton.edu/ mcdonald/examples/loopantenna.pdf.
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